

THE DETECTION OF LOW Eu ABUNDANCES IN EXTREMELY METAL-POOR STARS AND THE ORIGIN OF r -PROCESS ELEMENTS¹

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ABSTRACT

We report abundance analyses of three extremely metal-poor stars with $[\text{Fe}/\text{H}] \lesssim -3$ using the Subaru High Dispersion Spectrograph. All are found to have subsolar values of $[\text{Eu}/\text{Fe}]$. Comparison with our chemical evolution model of the Galactic halo implies the dominant source of Eu to be the low-mass end of the supernova mass range. Future studies of stars with low Eu abundances will be important to determine the r -process site.

Subject headings: Galaxy: evolution — Galaxy: halo — nuclear reactions, nucleosynthesis, abundances — stars: abundances — stars: Population II — supernovae: general

1. INTRODUCTION

The origin of the rapid neutron capture (r -process) elements is still uncertain. Although a few scenarios show some promise, such as neutrino winds (Woosley et al. 1994) in core-collapse supernovae (SNe), the collapse of ONeMg cores resulting from 8 to 10 M_{\odot} stars (Wanajo et al. 2003), and neutron star mergers (Freiburghaus, Rosswog, & Thielemann 1999), no consensus has been achieved.

Analyses have shown a large scatter in the abundances of neutron-capture elements such as Sr, Ba, and Eu among metal-poor stars in the Galactic halo (e.g., McWilliam et al. 1995; Ryan, Norris, & Beers 1996). If the scatter accurately portrays the state of the interstellar medium (ISM) from which they formed, large dispersions in excess of observational errors indicate that the ISM was not fully mixed at the early epoch of Galactic history and that metal-poor stars contain products from only one or a few SNe (Ryan, Norris, & Bessell 1991). The scatter possibly reflects variation in the yields of SNe from different mass progenitors; the huge dispersions in neutron-capture elements imply that r -process yields are highly dependent on the masses of SN progenitors.

Eu traces the r -process. The chemical evolution model of Ishimaru & Wanajo (1999, hereafter IW99) explains the large dispersion of Eu/Fe observed in very metal-poor stars if the r -process elements originate from a limited mass range of progenitor stars, such as the SNe from stars of 8–10 M_{\odot} or, alternatively, greater than 30 M_{\odot} (see also Travaglio et al. 1999). Tsujimoto, Shigeyama, & Yoshii (2000) have concluded by a similar approach that the SNe from $\approx 20 M_{\odot}$ stars are the dominant source of the r -process elements. A clear difference among these cases is expected in the stellar distribution of $[\text{Eu}/\text{Fe}]$ at $[\text{Fe}/\text{H}] \lesssim -3$ (IW99). However, a shortage of observational data in this metallicity range has made it difficult to distinguish between the proposed r -process sites.

¹ Based on data collected at Subaru Telescope, which is operated by the National Astronomical Observatory of Japan.

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⁶ $[X_i/X_j] \equiv \log(N_i/N_j) - \log(N_i/N_j)_{\odot}$, where N_i indicates abundance of i th element X_i .

In this Letter, we report on three extremely metal-poor stars that we show to have very low Eu abundances (§ 2). These data are compared with our chemical evolution models to distinguish between the proposed r -process sites (§ 3).

2. OBSERVATIONS AND ANALYSIS

We selected three very metal-poor ($[\text{Fe}/\text{H}] \lesssim -3$) giants—HD 4306, CS 22878-101, and CS 22950-046—which were known from previous studies (McWilliam et al. 1995; McWilliam 1998) to have $[\text{Ba}/\text{Fe}] \sim -1$, typical for their metallicities. Observations were made with the High Dispersion Spectrograph (Noguchi et al. 2002) of the 8.2 m Subaru Telescope in 2001 July, at a resolving power of $R = 50,000$. The total exposure time was 50, 240, and 360 minutes for HD 4306, CS 22878-101, and CS 22950-046, respectively, resulting in signal-to-noise ratios per 0.012 Å pixel at 4100 Å of 260, 110, and 80. Data reduction was performed in the standard way with IRAF.⁷

Equivalent widths were measured by fitting Gaussian profiles to the absorption lines of species listed in Table 1,⁸ except for Eu. An LTE analysis using model atmospheres of Kurucz (1993), based on equivalent widths measured above, was performed for species other than Eu. Results are given in Table 1. For Ba we included the hyperfine splitting and isotope shifts (McWilliam 1998), assuming isotope ratios of the r -process component in solar system material (Arlandini et al. 1999).

The atmospheric parameters (effective temperature, T_{eff} (K), gravity, g (cm s^{-2}), microturbulent velocity, v_{mic} (km s^{-1}), and metallicity that is assumed to be $[\text{Fe}/\text{H}]$ of CS 22878-101 and CS 22950-046 were adopted from Cohen et al. (2002) and Carretta et al. (2002); i.e., $T_{\text{eff}}/\log g/[\text{Fe}/\text{H}]/v_{\text{mic}} = 4775/1.3/-3.1/2.0$ and $4730/1.3/-3.3/2.0$, respectively. We found that these gravities satisfy the ionization balance between Fe I and Fe II, and the microturbulence velocities lead to no dependence of the derived abundance on equivalent widths of Fe I lines. The derived Fe abundances also agree very well with the assumed metallicities. So we adopted these atmospheric parameters without modification. For HD 4306, we adopted $T_{\text{eff}} = 5000$ K based on $V-K = 2.15$ (Norris, Bessell, & Pickles 1988; and Two Micron All Sky Survey data, Skrutskie et

⁷ IRAF is distributed by the National Optical Astronomy Observatories, which is operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.

⁸ The line list and measurements are available at <http://optik2.mtk.nao.ac.jp/~waoki/EW>.

TABLE 1
ELEMENTAL ABUNDANCES

Element	HD 4306				CS 22878-101				CS 22950-046			
	[X/Fe] ^a	<i>n</i>	σ^b	$\log \epsilon^c$	[X/Fe] ^a	<i>n</i>	σ^b	$\log \epsilon^c$	[X/Fe] ^a	<i>n</i>	σ^b	$\log \epsilon^c$
Fe I	-2.76	168	0.13	+4.74	-3.14	161	0.14	+4.36	-3.34	141	0.18	+4.16
Fe II	-2.79	19	0.13	+4.71	-3.07	20	0.14	+4.43	-3.34	17	0.18	+4.16
Sr II	+0.28	2	0.27	+0.42	-0.15	2	0.28	-0.33	-0.18	2	0.30	-0.60
Ba II	-1.09	2	0.14	-1.65	-0.73	2	0.16	-1.61	-1.31	2	0.19	-2.43
Eu II	-0.57	2	0.16	-2.80	-0.30	2	0.24	-2.85	<-0.2	1	...	<-3.0

^a [Fe/H] for Fe I and Fe II.

^b Uncertainty in [Fe/H] or [X/Fe] values.

^c $\log \epsilon(X) = \log_{10}(N_X/X_H) + 12.0$ for elements X.

al. 1997) and the temperature scale of Houdashelt, Bell, & Sweigart (2000). We derived g and v_{mic} from the analysis by requiring ionization balance and no dependence of abundance on line strength. The resulting parameters— $T_{\text{eff}}/\log g/[\text{Fe}/\text{H}]/v_{\text{mic}} = 5000/2.0/-3.0/1.85$ —are similar to those of other works (McWilliam et al. 1995; Honda et al. 2003).

For the three Eu II lines, at 3819, 4129, and 4205 Å, a spectrum synthesis technique was applied. The line list includes hyperfine splitting and isotope shifts from Lawler et al. (2001). Contamination by other species was included using the com-

prehensive line list of Kurucz & Bell (1995). Figure 1 compares the observed and synthetic spectra for the 3819 Å line, which is the strongest of the three. Eu is clearly detected in the spectra of HD 4306 and CS 22878-101, but not in CS 22950-046. We derived consistent results from the 4129 Å line, but the measurement of 4205 Å is uncertain because of severe contamination by V II at 4205.09 Å. We adopted the average of the Eu abundances derived from the 3819 and 4129 Å lines. In Table 2, we give the equivalent widths of the Eu II lines estimated from the spectrum synthesis using the Eu abundances derived from individual lines. For CS 22950-046, we estimated the upper limit of its equivalent widths from the 3 σ depth of the random noise level and typical line widths for this object, and determined the upper limit of the Eu abundance, given in Table 1. The continuum fit is constrained over a wider spectral range, which is shown in Figure 1; a systematic error of 1% would change the [Eu/H] limit from the 3819 Å line by 0.1.

The random error for each species is estimated from the standard deviation of the abundances derived from individual lines. The values are sometimes unrealistically small when the number of lines is small, so we adopted the larger of the values for the species and that for Fe I, 0.08–0.15 dex, depending on the species and the quality of the spectra. We estimated the errors in the Eu abundances by eye by fitting synthetic spectra to the observed ones, adopting 0.1 dex for HD 4306 and 0.2 dex for CS 22878-101. The errors in the abundance determinations due to uncertainties of the atmospheric parameters were evaluated for $\Delta T_{\text{eff}} = 100$ K, $\Delta \log g = 0.3$, $\Delta v_{\text{mic}} = 0.3$ km s⁻¹, and $\Delta[M/H] = 0.3$ for HD 4306, and were added in quadrature with the random error estimated above to give the total uncertainty in Table 1. The larger errors in [Sr/Fe] reflect the sensitivity to microturbulence of the two resonance lines. The [Eu/Fe] values are affected most significantly by the uncertainty in the surface gravity, by 0.12 for $\Delta \log g = 0.3$.

Comparisons between our results for the two CS objects and those by Carretta et al. (2002) show good agreement (within 0.17 dex) for the species in Table 1 other than Eu. An exception is Ba in CS 22878-101: our abundance is 0.35 dex higher than that by Carretta et al. (2002). The reason for this discrepancy is unclear. The equivalent widths of the Ba II $\lambda 4554$ line agree

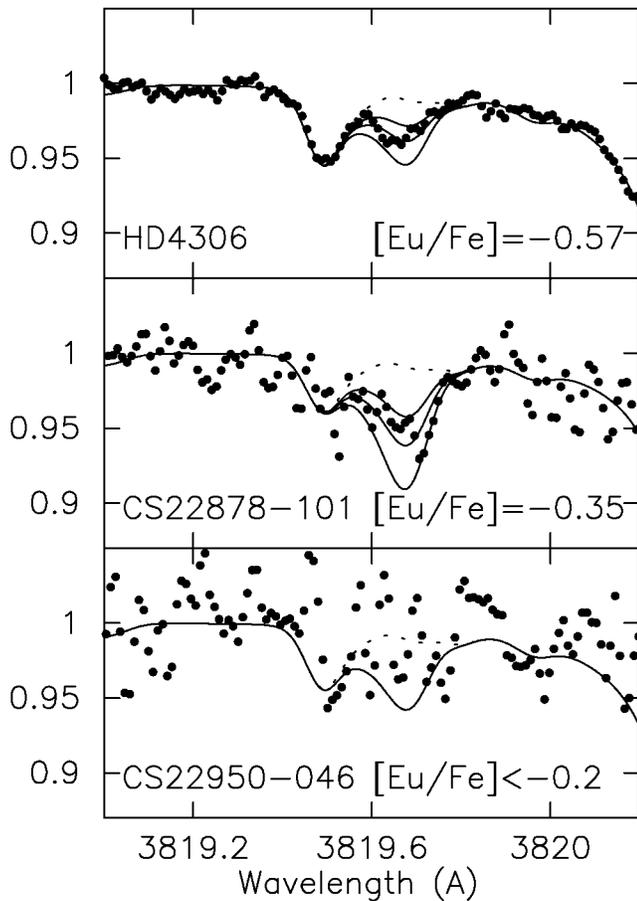


FIG. 1.—Comparison of the observed spectra (filled circles) and synthetic ones (lines) near the Eu II 3819.7 Å line. For HD 4306 and CS 22878-101, three synthetic spectra differ in step of $\Delta[\text{Eu}/\text{Fe}] = 0.2$ dex and are shown by solid lines. Central line is calculated for the Eu abundance presented in each panel. The solid line in the panel for CS 22950-046 indicates the synthetic spectrum for the upper limit of Eu abundance presented in the panel (see text). Dotted lines are the spectra calculated assuming no Eu contribution.

TABLE 2
EQUIVALENT WIDTH OF EU LINES

Wavelength (Å)	Equivalent Width (mÅ)		
	HD 4306	CS 22878-101	CS 22950-046
3819.7	3.8	7.4	<6.8
4129.7	2.2	5.3	...
4205.0	2.5	5.0	<3.4

well between the two studies, and the atmospheric parameters are the same. The difference in the assumed isotope ratios of Ba in the abundance analyses does not explain this discrepancy.

3. COMPARISON WITH THE GALACTIC CHEMICAL EVOLUTION MODELS

Assuming that star formation is induced by a single SN explosion, IW99 have constructed an inhomogeneous chemical evolution model of the Galactic halo. The chemical composition of a newly formed star is a mixture of the products from each SN and the ISM swept up by the expanding ejecta. This model, along with some improvements (see Ishimaru, Prantzos, & Wanajo 2003), leads to a large scatter of relative abundances for elements that have strong dependencies on the progenitor masses. The scattering of relative abundances in α and iron-peak elements, on the other hand, are predicted to be small (as is observed), owing to the weak dependencies of their yield ratios on the progenitor masses. We use this model to derive the Galactic evolution of the r -process tracer Eu. We consider three cases in which the r -process elements originate from SNe of (A) 8–10 M_{\odot} , (B) 20–25 M_{\odot} , and (C) greater than 30 M_{\odot} stars.

Yields for Type II and Type Ia SNe are taken from Nomoto et al. (1997a, 1997b). The 8–10 M_{\odot} stars are assumed to produce no iron, since their contribution to the enrichment of iron-peak elements in the Galaxy is negligible (Wanajo et al. 2003). The mass of Eu produced per SN event is assumed to be constant over the mass range of each case; this simplification will not affect the result significantly because of the narrow mass range of Eu-generating SNe assumed in each case. The requirement that the model reproduce the solar values $[\text{Eu}/\text{Fe}] = [\text{Fe}/\text{H}] = 0$ implies ejected Eu masses 3.1×10^{-7} , 1.1×10^{-6} , and $7.8 \times 10^{-7} M_{\odot}$ for cases A–C, respectively. These are reasonable amounts from a nucleosynthetic point of view (Wanajo et al. 2001, 2003).

Figure 2 compares the results of these cases with the observational data from § 2 (*open circles*) and from sources cited in the legend (*filled circles*). The predicted log number density of stars per unit area on the $[\text{Fe}/\text{H}]$ - $[\text{Eu}/\text{Fe}]$ plane is also shown as a color image.

The heavy solid lines indicate the average values of $[\text{Eu}/\text{Fe}]$ as a function of metallicity, calculated from the models. The medium and light solid lines indicate, respectively, the 50% and 90% confidence lines.

Figure 2 shows that observable differences between the cases appear at $[\text{Fe}/\text{H}] \lesssim -3$. In case C (Fig. 2c), most of the stars are expected to have $[\text{Eu}/\text{Fe}] > 0$, owing to Eu production solely by massive, short-lived stars. In cases A (Fig. 2a) and B (Fig. 2b), significant numbers of stars having $[\text{Eu}/\text{Fe}] < 0$ are predicted at $[\text{Fe}/\text{H}] \lesssim -3$, owing to the delayed production of Eu by lower mass SN progenitors. This is outstanding in case A, which has the largest delay.

Most previous observational data, many of them upper limits, distribute between the 90% confidence lines for all cases, which has made it difficult to determine the mass range of the r -process site. Our newly obtained data add the lowest detections of Eu, at $[\text{Fe}/\text{H}] \lesssim -3$, and help distinguish between the three cases. The best agreement can be seen in case A (Fig. 2a), in which the three stars, and most other stars from previous observations, are located between the 50% confidence lines at $[\text{Fe}/\text{H}] \lesssim -3$. In case B (Fig. 2b), all three stars are located below the average (*heavy solid line*), although they

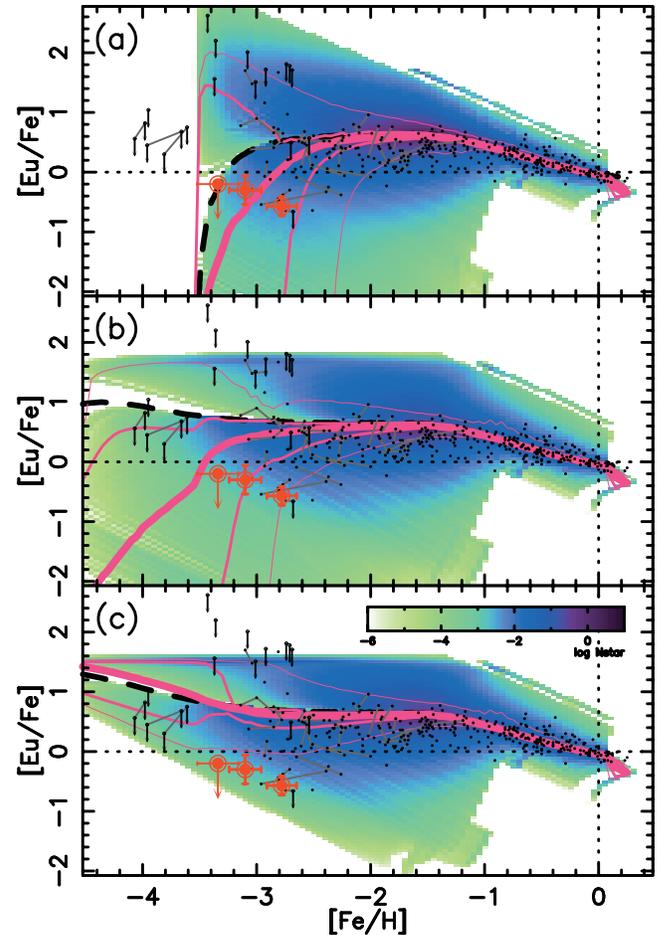


FIG. 2.—Comparison of the observed data with the model predictions. The r -process site is assumed to be SNe of (a) 8–10 M_{\odot} , (b) 20–25 M_{\odot} , and (c) greater than 30 M_{\odot} stars. The predicted number density of stars per unit area is color coded. The average stellar abundance distributions are indicated by heavy lines, with 50% and 90% confidence intervals (*medium and light lines, respectively*). The average abundances of the ISM are denoted by the heavy dashed lines. The current observational data are given by open circles, with others (*filled circles*) taken from Gratton & Sneden (1994), McWilliam et al. (1995), McWilliam (1998), Woolf, Tomkin, & Lambert (1995), Ryan et al. (1996), Shetrone (1996), Sneden et al. (1996), Westin et al. (1998), Burris et al. (2000), Fulbright (2000), Norris, Ryan, & Beers (2001), Johnson (2002), Johnson & Bolte (2002), François et al. (2003), and Honda et al. 2003).

remain above the 90% confidence line. In case C (Fig. 2c), these stars are located outside the 90% confidence region. We suggest, therefore, that case A is most likely to be the r -process site (i.e., SNe from low-mass progenitors such as 8–10 M_{\odot} stars). More detections of Eu at low metallicity will be needed to clearly distinguish between cases A and B. Note that this result is not due to a selection bias, since the $[\text{Ba}/\text{Fe}]$ values (~ -1) are typical of extremely metal-poor stars with $[\text{Fe}/\text{H}] \sim -3$ (Ryan et al. 1996).

Our analysis gives $[\text{Ba}/\text{Eu}]$ values consistent with the solar r -process (Arlandini et al. 1999) when estimated errors are included (see Table 3). Hence, our result may hold for heavy r -process elements with $Z \geq 56$, not just $Z \approx 63$. The values of $[\text{Sr}/\text{Ba}]$, however, are significantly higher than for the solar r -process, implying that these three stars exhibit light r -process elements ($Z < 56$) produced in more massive SNe (greater than 10 M_{\odot}).

The discussion above suggests that the production of the

TABLE 3
ABUNDANCE RATIOS AMONG Sr, Ba, AND Eu

Ratio	HD 4306	CS 22878-101	CS 22950-046	Solar ^a
[Sr/Ba]	1.37 ± 0.30	0.58 ± 0.32	1.13 ± 0.36	-0.10
[Ba/Eu]	-0.52 ± 0.21	-0.43 ± 0.29	>-1.11	-0.69

^a Solar system r -process abundances by Arlandini et al. 1999.

r -process elements is associated with a small fraction of SNe near the low-mass end of the range. Neutrino winds in the explosions of massive stars may face difficulties as a dominant source of the r -process elements. Wanajo et al. (2001) have demonstrated that an r -process in the neutrino winds proceeds from only very massive proton-neutron stars, which might result from massive progenitors such as ≥ 20 – $30 M_{\odot}$ stars, which is similar to case B. Hypernovae (>20 – $25 M_{\odot}$; Maeda & Nomoto 2003) or pair-instability SNe (140 – $260 M_{\odot}$; Heger & Woosley 2002) resulting from stars near the high-mass end of the SN progenitors, similar to case C, are clearly excluded as major r -process sites.

We suggest, therefore, that the dominant r -process site is SN

explosions of collapsing ONeMg cores from 8 – $10 M_{\odot}$ stars (Wheeler, Cowan, & Hillebrandt 1998). Recently, Wanajo et al. (2003) have demonstrated that the prompt explosion of the collapsing ONeMg core from a $9 M_{\odot}$ star reproduces the solar r -process pattern for nuclei with $A > 130$, and is characterized by a lack of α elements and only a small amount of iron-peak elements. This clearly differs from more massive SNe with iron cores (greater than $10 M_{\odot}$) that eject both these elements, and is consistent with the fact that the abundances of heavy r -process elements in stars with $[\text{Fe}/\text{H}] \sim -3$ are not related to those of iron-peak elements or of elements with lower atomic numbers (Qian & Wasserburg 2003).

This study shows the importance of detecting Eu in extremely metal-poor stars to explore the origin of r -process elements. Further observations are needed to confirm the origin.

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