

SGR J1550 – 5418 BURSTS DETECTED WITH THE *FERMI* GAMMA-RAY BURST MONITOR DURING ITS MOST PROLIFIC ACTIVITY

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ABSTRACT

We have performed detailed temporal and time-integrated spectral analysis of 286 bursts from SGR J1550 – 5418 detected with the *Fermi* Gamma-ray Burst Monitor (GBM) in January 2009, resulting in the largest uniform sample of temporal and spectral properties of SGR J1550 – 5418 bursts. We have used the combination of broadband and high time-resolution data provided with GBM to perform statistical studies for the source properties. We determine the durations, emission times, duty cycles and rise times for all bursts, and find that they are typical of SGR bursts. We explore various models in our spectral analysis, and conclude that the spectra of SGR J1550 – 5418 bursts in the 8 – 200 keV band are equally well described by optically thin thermal bremsstrahlung (OTTB), a power law with an exponential cutoff (Comptonized model), and two black-body functions (BB+BB). In the spectral fits with the Comptonized model we find a mean power-law index of -0.92 , close to the OTTB index of -1 . We show that there is an anti-correlation between the Comptonized E_{peak} and the burst fluence and average flux. For the BB+BB fits we find that the fluences and emission areas of the two blackbody functions are correlated. The low-temperature BB has an emission area comparable to the neutron star surface area, independent of the temperature, while the high-temperature blackbody has a much smaller area and shows an anti-correlation between emission area and temperature. We compare the properties of these bursts with bursts observed from other SGR sources during extreme activations, and discuss the implications of our results in the context of magnetar burst models.

Subject headings: pulsars: individual (SGR J1550 – 5418; 1E 1547.0-5408; PSR 1550 – 5418) – stars: neutron – X-rays: bursts

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1. INTRODUCTION

Magnetars are members of the diverse neutron star family: they are isolated neutron stars whose persistent X-ray emission and soft gamma-ray bursts are powered by their extremely strong magnetic fields of 10^{14} – 10^{15} G (Duncan & Thompson 1992; Kouveliotou et al. 1998; Thompson et al. 2002). Historically two classes of sources have been identified with characteristics that could be explained by the magnetar model: Anomalous X-ray Pulsars (AXPs) and Soft Gamma Repeaters (SGRs). These original classifications were based on several distinct differences between the two populations (for reviews, see Woods & Thompson 2006; Mereghetti 2008), but it has now become clear that they share several properties. Therefore, AXPs and SGRs form one source class with small differences still maintained, most notably the rates and energetics of bursts during active episodes, with the luminosity distributions of AXP bursts being lower by orders of magnitude than those of SGR bursts.

The source discussed in this paper was discovered by the *Einstein* X-ray satellite and named 1E 1547.0-5408 (Lamb & Markert 1981); it was subsequently confirmed in the *ASCA* Galactic plane survey (Sugizaki et al. 2001). *Chandra* and *XMM-Newton* observations identified the source as a magnetar candidate based on its varying X-ray flux, spectral characteristics and the tentative

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association with the supernova remnant G327.24-0.13 (Gelfand & Gaensler 2007), although no X-ray pulsations or soft gamma-ray bursts were detected. The magnetar nature of the source was confirmed by the discovery of radio pulsations with the Parkes radio telescope, which were subsequently associated with 1E1547.0-5408 with ATCA observations; the radio source was named PSR J1550-5418 (Camilo et al. 2007). It is one of only three magnetars exhibiting pulsed radio emission, the other ones being XTE J1810-197 (Camilo et al. 2006) and PSR J1622-4950 (Levin et al. 2010). A surface dipole magnetic field strength of 2.2×10^{14} G has been inferred from its spin period of 2.07 s (the shortest of all magnetars) and its period derivative of 2.32×10^{-11} s s⁻¹. Further studies of the radio and persistent X-ray emission showed that the rotation and magnetic axes of this neutron star are nearly aligned (Halpern et al. 2008; Camilo et al. 2008).

Given its characteristics and the lack of bursting behavior, the magnetar source was originally classified as an AXP (Camilo et al. 2007). However, in 2008 the source started a series of three active episodes that covered over half a year: the first was in October 2008, the second, most burst active one started in January 2009, and the last in March 2009. During the first episode, the *Swift* Burst Alert Telescope (BAT; Krimm et al. 2008a,b; Israel et al. 2010) and the *Fermi* Gamma-ray Burst Monitor (GBM; von Kienlin & Briggs 2008; van der Horst & Briggs 2008) detected several tens of bursts; as a result the source was reclassified as an SGR and designated SGR J1550 – 5418. This reclassification was firmly established by the extreme bursting activity displayed on January 22, 2009, when hundreds of bursts were detected in one day by several instruments, i.e., *Swift* (Gronwall et al. 2009; Scholz & Kaspi 2011), *Fermi*/GBM (Connaughton & Briggs 2009; von Kienlin & Connaughton 2009), *INTEGRAL*/SPI-ACS (Savchenko et al. 2009; Mereghetti et al. 2009a,b) and IBIS/ISGRI (Savchenko et al. 2010), *Suzaku* WAM (Terada et al. 2009), RHESSI (Bellm et al. 2009), and KONUS-WIND (Golenetskii et al. 2009a,b,c). These bursts had durations and luminosities typical of SGR bursts; in fact some were so bright that they caused ionospheric disturbances observed in very low frequency radio wave data (Chakrabarti et al. 2009; Tanaka et al. 2010).

The distance to the source is currently uncertain, varying from ~ 9 kpc, based on the radio dispersion measure (Camilo et al. 2007), down to $\sim 4 - 5$ kpc, based on the tentative supernova remnant association (Gelfand & Gaensler 2007) and the detection of expanding dust-scattering X-ray rings (Tiengo et al. 2010). The former estimate has large uncertainties arising from the Galactic free electron density model of Cordes & Lazio (2002) used for calculating the radio dispersion measure. Therefore, we adopt here a source distance of 5 kpc. Throughout this paper we use the notation $d_5 = \text{distance}/(5 \text{ kpc})$, when calculating burst energies and luminosities.

In this paper we present the temporal and time-integrated spectral analysis of 286 SGR J1550 – 5418 bursts detected with GBM in January 2009 for which we have high time-resolution data. Given the very large field of view (the entire un-occulted sky) of GBM, this is the largest uni-

form sample of temporal and spectral properties of SGR J1550 – 5418 bursts during the January activation. This data set enables us to determine their durations, spectral shapes and energetics, and compare them with other magnetar sources during similar high bursting activity periods: SGR 1900+14 (Gögüş et al. 1999; Israel et al. 2008), SGR 1806-20 (Gögüş et al. 2000), AXP 1E2259+586 (Gavriil et al. 2004), and SGR 1627-41 (Esposito et al. 2008). We also compare our results with the analysis of bursts detected by *Swift* during the same time period (Scholz & Kaspi 2011), and discuss them in the framework of current theoretical models. Here we focus on the January 2009 bursts, while the analysis of the October 2008 and February - March 2009 bursts are discussed in von Kienlin et al. (in preparation) and Younes et al. (in preparation). The discovery of enhanced persistent emission in the GBM data during the beginning of activity on January 22 has been presented in Kaneko et al. (2010).

In Section 2 we present our data selection of GBM bursts, and in Sections 3 and 4 we present our temporal and spectral analysis, respectively. We discuss our results and their interpretation in the context of other bursting magnetar sources in Section 5, and we summarize in Section 6.

2. SAMPLE AND DATA SELECTION

GBM comprises 14 detectors with a field of view of 8 sr (the entire sky un-occulted by the Earth): twelve NaI detectors covering a spectral range from 8 keV to 1 MeV, and two BGO detectors covering 0.2 – 40 MeV (see Meegan et al. 2009, for an overview of the instrument and its capabilities). We use here only data from the NaI detectors because SGR bursts are not detected above 200 keV.

GBM collects three data types: Time-Tagged Event (TTE), CTIME and CSPEC (Meegan et al. 2009). To perform detailed temporal and spectral analyses over the typical burst durations of ~ 0.1 s, we selected for our sample only events with TTE data, which provide a time-tagged photon event list with a temporal resolution as low as $2 \mu\text{s}$ in 128 energy channels logarithmically spaced over the energy range of the NaI detectors. The other two data types cannot be used for our purposes, because either their accumulation times are much longer (CSPEC) or they have a worse spectral resolution (CTIME). The TTE data type covers the time range from 30 seconds prior to 300 seconds after a trigger, and GBM has been designed to not trigger for 600 seconds after a previous trigger. This leaves at least 270 s between consecutive triggers without TTE data. Fortunately, SGR J1550 – 5418 triggered GBM very frequently (41 times on January 22 alone), providing a good TTE data coverage for hundreds of bursts.

To perform a comprehensive and systematic study of the SGR J1550 – 5418 bursts, we applied our untriggered burst search algorithm to the CTIME data type in both continuous (0.256 s time resolution) and trigger (0.064 s time resolution) modes (Kaneko et al. 2010). For each burst we required a count rate above background in the 10 – 50 keV energy range of at least 5.5σ and 4.5σ in the first and second brightest detectors, respectively. Note that the GBM trigger criteria differ from the ones applied in our untriggered burst search, resulting in omit-

ting several GBM triggers from our sample. We then compared each burst candidate light curve and spectral behavior to those of typical SGR bursts, and checked if the SGR J1550 – 5418 location was consistent with the relative rates in the NaI detectors at the time of the burst (for the orientation of the spacecraft at that time). The untriggered burst search resulted in 555 events on January 22 alone and 597 events in total up to January 29. This total number, however, includes events for which TTE data are not available and thus excluded from this analysis. Moreover, in several occasions multiple events were part of the same burst as defined by the method employed by Gögüş et al. (2001). More specifically, we required that for two events to qualify as two separate bursts, the time between their peaks in the TTE data had to be greater than a quarter of the spin period of the SGR (0.5 s in the case of SGR J1550 – 5418), and the count rate level had to drop to the background level between the peaks. Applying these criteria to our initial list of events, we collected 286 bursts from January 22 to 29, which is the sample we used for our detailed temporal and spectral analysis.

There is a caveat, however, associated with our sample completion. For a significant fraction of the time we were not able to detect SGR J1550 – 5418 with GBM: during *Fermi* passages through the South Atlantic Anomaly (SAA) and when the source was occulted by the Earth. Specifically, on January 22 *Fermi* was in the SAA for 3.5 hours and SGR J1550 – 5418 was occulted by the Earth for almost 8 hours. Since some of these times overlap, the total time that GBM could not detect bursts from SGR J1550 – 5418 was ~ 11 hours. To obtain a complete picture of the source activity, we searched the data of other satellites that were triggered by the source. The only instrument with a continuous SGR J1550 – 5418 coverage during the entire day of 2009 January 22 was the *INTEGRAL*/SPI-ACS. The SPI-ACS, however, has a higher energy threshold (>80 keV) and provides no energy resolution. As a result it detected fewer bursts (slightly over 200 Mereghetti et al. 2009b), of which only 84 were also detected by the *INTEGRAL*/IBIS/ISGRI instrument, and for which spectral information (> 20 keV) was available (Savchenko et al. 2010). Thus the GBM sample is by far the most complete uniform data set for that entire day’s activity from SGR J1550 – 5418.

For each one of the 286 GBM bursts we used the TTE data of those NaI detectors that had a source viewing angle (angle between the source direction and the detector normal) of less than 60° . We then checked whether each one of these detectors had an unobstructed view of SGR J1550 – 5418, i.e., if no parts of the spacecraft or the Large Area Telescope (LAT) onboard *Fermi* were obstructing the view. This last check is important because several bursts had such a high peak flux that they caused Autonomous Repoint Requests of *Fermi*, leading to the re-pointing of the spacecraft to the source. The final orientation in those cases was 5° off from the LAT zenith, resulting in the LAT obstructing the view for some detectors.

Finally, several bursts in our sample were so bright that they saturated the High Speed Science Data Bus of GBM. This effect occurs when the total TTE count rate of all detectors exceeds a limit of $375,000$ counts s^{-1} . Out

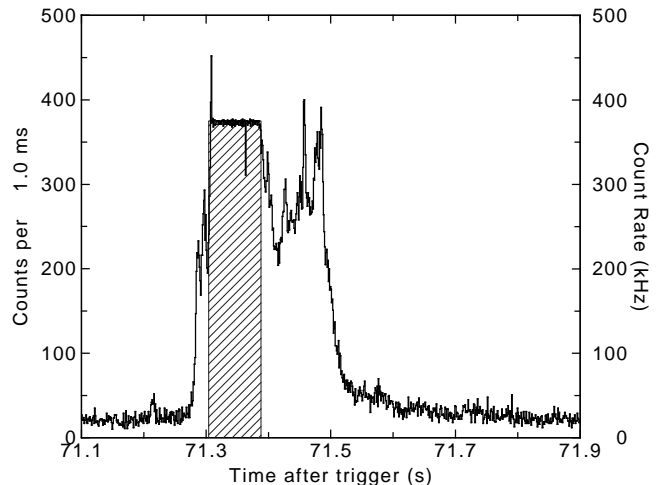


FIG. 1.— One of the 23 saturated SGR J1550 – 5418 bursts (06:59:34 UT on 22 January 2009). The light curve is the sum of the count rates in all 14 detectors across the entire energy range of GBM. The saturated time interval is indicated by the hatched area.

of the 286 bursts only 23 were affected; an example of a saturated burst is shown in Figure 1 (hatched area). We did not use the saturated bursts in our temporal analyses (Section 3), since the durations, emission times and peak times of these bursts are severely affected. We performed a spectral analysis using the non-saturated parts of these bursts, and used their fluences only in the compilation of the cumulative energy fluence emitted during the active period of SGR J1550 – 5418 (Section 4). None of the 23 saturated bursts were used in any of the other parts of the spectral analysis and discussion sections in this paper.

3. TEMPORAL ANALYSIS

To characterize the temporal properties of all the bursts in our sample, we determined their durations, emission times, duty cycles and rise times. We calculate these quantities using TTE data binned at 2 ms resolution in the 8–100 keV energy range. We apply a two-step background subtraction method to the total number of counts during and around the burst time interval; details on the particular method used here are given in Gögüş et al. (2001) and Lin et al. (2011).

The T_{90} (T_{50}) duration is the time during which 90% (50%) of the total burst counts are accumulated (Kouveliotou et al. 1993). Another parameter to quantify burst durations is the emission time τ_{90} (τ_{50}) (Mitrofanov et al. 1999). In contrast with $T_{90,50}$, the emission time does not necessarily span a time interval of consecutive time bins; it is computed by adding time bins ordered from the highest to the lowest number of counts until 90% or 50% of the burst counts is accumulated. We then calculated for each burst the duty cycle δ_{90} (δ_{50}), which is the ratio of τ_{90} (τ_{50}) over T_{90} (T_{50}); this quantity gives an indication of how sharply peaked a given burst light curve is (Mitrofanov et al. 1999). The final temporal parameter we determined is the burst rise time T_{rise} (the time for a burst to reach its peak count rate); in multiple peaked events, we used the highest count rate bin as the peak.

The results of our temporal analysis are shown in Fig-

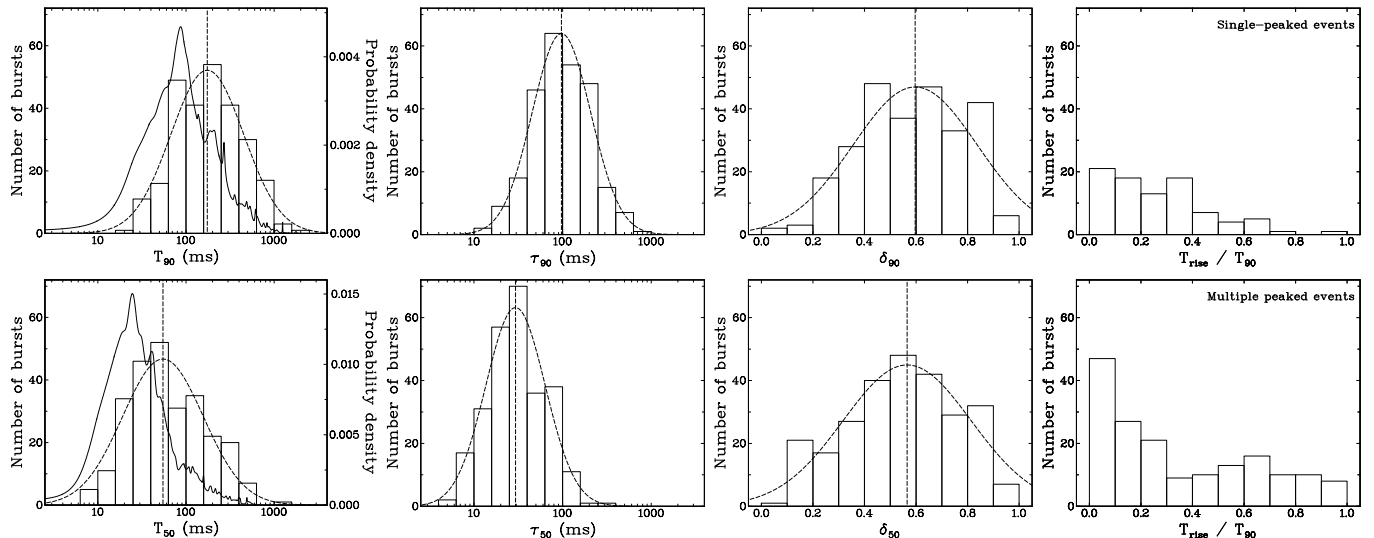


FIG. 2.— Distributions of the durations $T_{90,50}$ (first column panels); emission times $\tau_{90,50}$ (second column panels); duty cycles $\delta_{90,50}$ (third column panels); and T_{rise}/T_{90} (fourth column panels) for single-peaked (top) and multiple peaked (bottom) events. The dashed lines show the best-fit log-normal ($T_{90,50}$ and $\tau_{90,50}$) and normal ($\delta_{90,50}$) functions, with the vertical dashed lines indicating the fitted mean values. For $T_{90,50}$ the probability density functions are shown as solid lines.

TABLE 1
TEMPORAL ANALYSIS OF 263 SGR J1550 – 5418 BURSTS

Parameter	Fit mean ^a	σ^b	Mean
T_{90} (ms)	174 ± 10	0.41 ± 0.02	258
T_{50} (ms)	55 ± 4	0.46 ± 0.03	104
τ_{90} (ms)	97 ± 3	0.33 ± 0.01	127
τ_{50} (ms)	29 ± 2	0.34 ± 0.02	39
δ_{90}	0.60 ± 0.03	0.24 ± 0.04	0.58
δ_{50}	0.56 ± 0.03	0.25 ± 0.03	0.54

^alog-normal fit for $T_{90,50}$ and $\tau_{90,50}$, normal fit for $\delta_{90,50}$
^bin the log-frame except for $\delta_{90,50}$

ure 2 and Table 1. Figure 2 shows the distributions of $T_{90,50}$ (first column panels), $\tau_{90,50}$ (second column panels) and $\delta_{90,50}$ (third column panels), and T_{rise}/T_{90} for single-peaked and multiple peaked events (fourth column panels). We fit the distributions of the durations and emission times with log-normal functions, and the duty cycle distributions with normal functions. The means and standard deviations of these functions are given in Table 1, together with their normal mean values.

Next we used the $T_{90,50}$ values and their uncertainties to construct their probability density functions (PDFs) as described in Lin et al. (2011). For each individual PDF we adopted a two-sided normal distribution with the widths given by the measured (asymmetric) uncertainties (Starling et al. 2008). Figure 2 shows the average of all the individual PDFs for $T_{90,50}$. Note that the PDFs peak at smaller values than the duration histograms, namely at ~ 85 ms and ~ 25 ms for T_{90} and T_{50} , respectively. This is due to the fact that the shorter (and weaker) events have relatively larger asymmetries in their uncertainties.

The histograms shown in Figure 2 for the various temporal parameters are typical of magnetar bursts, although there may be variations between different sources (e.g., Göğüş et al. 2001; Gavriil et al. 2004; Lin et al. 2011). The mean T_{90} values in the GBM en-

ergy band typically range between 0.1–0.2 s, while τ_{90} ranges are 0.05–0.1 s, resulting in duty cycles of 0.4 – 0.6 (see also Lin et al. 2011). Our mean T_{90} values for SGR J1550 – 5418 are larger than those found for *INTEGRAL* (68 ms; Savchenko et al. 2010), but smaller than the *Swift* bursts (305 ms; Scholz & Kaspi 2011). This is due to the different sensitivities and energy ranges of the instruments (see also Scholz & Kaspi 2011). Regarding the rise times, we confirm what has been established for other magnetar bursts (e.g., Göğüş et al. 2001; Scholz & Kaspi 2011): the asymmetric T_{rise}/T_{90} distributions for single-peaked events indicate that magnetar bursts have shorter rise than decay times, and the bimodal distribution of multiple peaked events shows that for most bursts the first peak is the brightest. In Section 5 we elaborate further on the correlation of the temporal and spectral properties of SGR J1550 – 5418, and compare these to similar correlations in other magnetar sources.

4. SPECTRAL ANALYSIS

We performed time-integrated spectral analysis of all 286 bursts in our sample using the spectral analysis software package *RMFIT* (*3.3rc8*)²⁰, which was developed specifically for GBM data analysis. For every burst we selected intervals of background without bursts present before and after a burst, and fit them with a polynomial of the third or fourth order. For the spectral fits we generated Detector Response Matrices using *GBMRSP v1.81*. To obtain the best fit parameters for any given spectral model we minimized the Castor C-statistic (C-statistic hereafter). This method is used to fit data with a low number of counts and is a modification²¹ of the Cash statistic so that it asymptotically distributes as χ^2 .

We fit various spectral models to the TTE data of each burst: a power law (PL), a black-body function (BB), optically thin thermal bremsstrahlung (OTTB), a power law with an exponential cutoff (Comptonized model), a

²⁰ <http://fermi.gsfc.nasa.gov/ssc/data/analysis/user/>

²¹ heasarc.gsfc.nasa.gov/docs/xanadu/xspec/wstat.ps

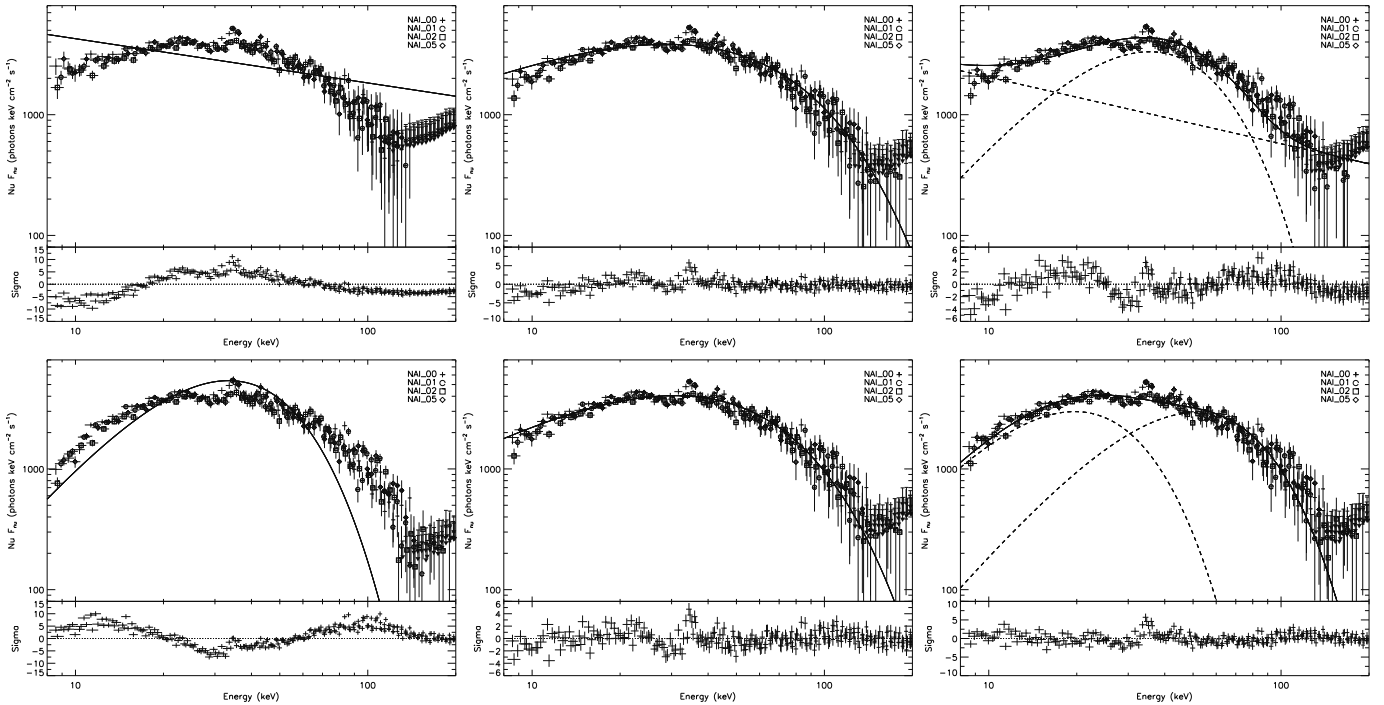


FIG. 3.— Time-integrated νF_ν spectra of one of the brightest SGR J1550 – 5418 bursts (06:59:34 UT on 22 January 2009), shown in Figure 1, for various spectral models: power law (top left), single black body (bottom left), OTTB (top center), Comptonized (bottom center), power law plus black body (top right), and two black body functions (bottom right). The bottom panel of each spectrum shows the fit residuals. The feature between 30 and 40 keV in the fit residuals is a result of the imperfect modeling of the NaI K-edge.

combination of a power law and a blackbody function (PL+BB), and two black-body functions (BB+BB). Figure 3 shows spectra of one of our brightest bursts (using only the unsaturated parts), whose light curve is shown in Figure 1. We have used in this spectral analysis counts from four NaI detectors following the detector selection criteria given in Section 2. The νF_ν spectra in Figure 3 show fits with the six different spectral models mentioned above and their fit residuals. The fit residual trends shown here are very similar for all bursts in our sample, with the brightest bursts presenting the strongest trends. From Figure 3 an obvious conclusion would be that OTTB, Comptonized and BB+BB fits provide a better description of the data than PL, BB or PL+BB. Section 4.1 attempts to quantify this statement. We note that the apparent feature between 30 and 40 keV in the fit residuals is a result of the NaI K-edge, which has not been modeled perfectly in the GBM calibration (Bissaldi et al. 2009). We have performed fits with and without the K-edge energy channels, and found that their inclusion does not change any of the fit parameters significantly, therefore we included these channels in all fits for better statistics.

4.1. Simulations

We have performed extensive simulations with *RMFIT* (3.4rc3) to evaluate the effect of statistical fluctuations on identifying the best model to describe SGR burst spectra. For this purpose we selected six bright bursts (two from SGR J1550 – 5418 and three from SGR J0501+4516; see also Lin et al. 2011), and compared pairs of various models, generating 30,000 synthetic spectra for each burst and each relevant detector. Each synthetic spectrum was based on the sum of the

predicted source counts and the measured background counts in each energy channel. The former counts were computed from the analytical function with the largest C-statistic used to fit the real data (null hypothesis) folded with the instrumental response function of the relevant detector (as a check, we also used as null hypothesis the function giving the lowest C-statistic and we obtained consistent results). The background counts were estimated for each detector from the real data. Next, we applied Poisson fluctuations to the summed counts to obtain the final synthetic spectrum.

During the fit process we generated a synthetic background spectrum by adding Poisson fluctuations to the real background estimate in each energy channel, and subtracting it from the synthetic spectrum. The resulting spectrum was fit with the two models of each pair. We checked and confirmed that the input spectral parameters were well recovered in the fits of the synthetic spectra, by comparing them and their statistical errors to the simulated parameter distributions and their uncertainties. Finally, to test whether this recovery was intensity dependent, we also simulated two weak events and arrived at the same conclusions.

We then proceeded to compare the fits between different model pairs. To this end we computed the C-statistic difference between one input model and each other spectral model for all 30,000 synthetic spectra per burst, and compared the resulting distribution to the difference obtained using the same models with the real burst data.

First we compared PL and BB with all other spectral models, because the former systematically gave the largest fit residual patterns for the majority of the events in our sample. For all simulated events, we found that statistical fluctuations in the signal and background can-

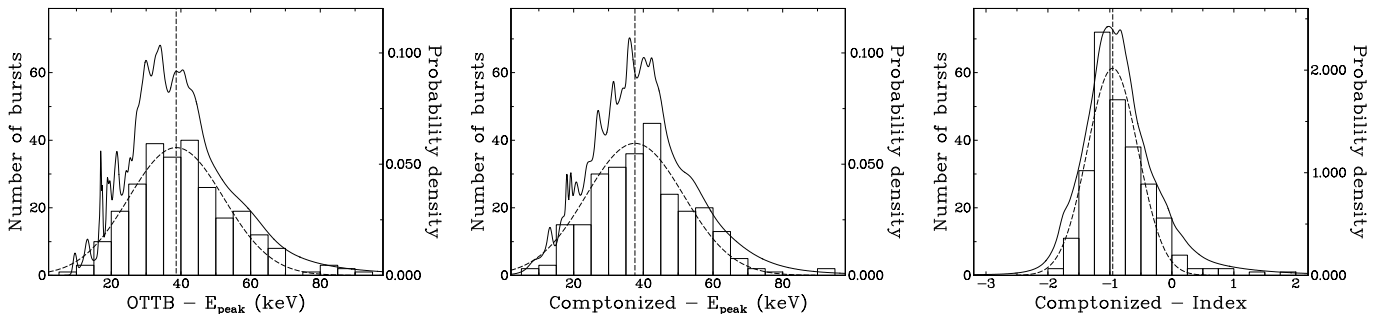


FIG. 4.— Distributions of the OTTB E_{peak} (left), Comptonized E_{peak} (middle), and of the Comptonized power-law index (right). The dashed lines show the best-fit normal functions, with the vertical dashed lines indicating the fitted mean values. The probability density functions are shown as solid lines.

not explain that OTTB or any of the more complex models fit the data better than PL or BB, including the weakest one with a fluence of 6.7×10^{-8} erg/cm². Using the same procedure, we then compared different models, namely OTTB, Comptonized, PL+BB and BB+BB. Our simulations showed that the statistical fluctuations in the data could account for the difference in C-statistic values between these models, preventing us from drawing conclusions on the best spectral shape. We cross-checked and confirmed the results of our simulations using XSPEC (v12.6; Arnaud 1996). We conclude, therefore, that we cannot unambiguously determine whether OTTB, Comptonized, PL+BB or BB+BB is the best spectral fit model for our burst sample.

The same method has been used to distinguish between spectral models for gamma-ray bursts (GRBs) detected with GBM (see e.g. Guiriec et al. 2011), albeit resulting in identification of the best-fit spectral model. The discrepancy with the SGR simulation results here could be explained by the fact that even the brightest SGR bursts have a lower total number of source counts mostly distributed over a much narrower energy range (10–50 keV).

4.2. Spectral Model Selection

Although our simulation results were inconclusive regarding the preferred spectral model, we attempt here to narrow down the number of models that give a good description of the SGR burst spectra. We did this down-selection based on the following arguments. When fitting the spectra of individual bursts, we obtained systematically large fit residual patterns with the PL, BB, and PL+BB functions for a large fraction of the events in our sample; all other model residuals were randomly distributed. For the PL and BB models we showed in Section 4.1 that they give a worse description of the data than all other models for bursts with fluences above 6.7×10^{-8} erg/cm². In addition, the PL and BB models had predominantly the largest C-statistic values for all bursts compared to all other models, resulting in very large C-statistic differences, typically tens to hundreds and for some events even a few thousand. This difference increased with fluence, indicating that the brightest events are worst fit with a PL or BB. Such a dependence on brightness was not seen when comparing C-statistic differences between all other models. The combination of all these arguments led us to discard PL and BB models from further analysis for bursts at all fluences. Our PL+BB fits gave a C-statistic value better than

OTTB for only one third of the bursts in our sample, while PL+BB has two more free parameters than OTTB. Moreover, the PL+BB model systematically gave residual patterns not seen in the BB+BB, Comptonized and OTTB fits. We, therefore, excluded the former model from further consideration.

It is not possible to exclude any of the three remaining models, namely OTTB, Comptonized and BB+BB. We cannot choose between OTTB and Comptonized because the power-law index we retrieve in the Comptonized fits is often very close to -1 (in photon space) and OTTB is basically a special case of Comptonized with an index of -1 . This is illustrated in the rightmost panel of Figure 4, which shows the distribution of the Comptonized power-law indices for all events, together with a normal fit to the distribution with a mean value of -0.95 and a width of 0.41 . Finally, BB+BB is very similar to the Comptonized function over this relatively small energy range (e.g., see the middle right and bottom right panels of Figure 3). When comparing the C-statistic for individual bursts, the difference between these two models lies almost always between -25 and 25 , with the majority between -10 and 10 , making the choice of the best fit very difficult. We, therefore, describe below our results of the spectral fits on all 263 events using the OTTB, Comptonized and BB+BB models.

4.3. OTTB and Comptonized Model

The first two panels in Figure 4 show the E_{peak} distributions for the OTTB (left panel) and the Comptonized (middle panel) models. We note that these are very similar with mean E_{peak} values of 39 keV and widths of 13 keV in both cases. The OTTB function has been applied successfully to SGR bursts detected with instruments that have a similar or smaller spectral coverage to GBM (e.g. Göğüş et al. 1999, 2000, 2001). Several studies, however, have shown that if the spectral range is extended down to several keV, OTTB overestimates the flux at lower energies (e.g., Fenimore et al. 1994; Feroci et al. 2004). Figure 3 indicates a small trend for overestimation below 10 keV, and a similar hint is seen in several other bright bursts in our sample. Figure 3 also shows that a Comptonized function with an index of ~ -0.7 corrects for this possible trend. For the brightest bursts pulse pile-up could affect the burst hardness, which would be reflected in the index value. We have performed simulations to study this effect, and for the count rates observed in our brightest bursts, the Comptonized index becomes flatter by at most 0.1 , while the

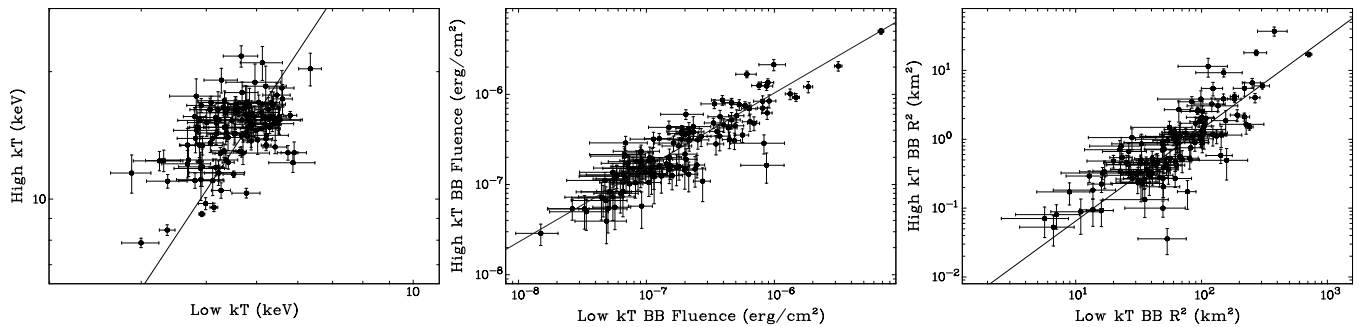


FIG. 5.— Correlations between the cool and hot blackbody parameters: temperatures (left panel), fluences (middle panel), and emission areas (right panel).

E_{peak} value does not change significantly. Given that the GBM calibration (Bissaldi et al. 2009) was done for energies starting at ~ 14 keV and then extrapolated to lower energies, at this point we cannot determine the significance of a Comptonized index deviation of 0.3 from -1 . We thus compare in Section 5 the spectral parameters derived with both models to other SGR sources.

4.4. Two blackbody Model

As an alternative to OTTB or Comptonized, several studies have shown that SGR bursts can be fit well with two blackbody functions (e.g. Feroci et al. 2004; Olive et al. 2004; Nakagawa et al. 2007; Israel et al. 2008; Esposito et al. 2008; Lin et al. 2011). We have fit 263 of our bursts in our sample (excluding the 23 saturated events) with BB+BB. Several of these, however, were faint and their parameters could not be well constrained. Therefore, in the following we have limited our sample to those bursts for which the temperatures and fluences of both black bodies have values that are at least two times larger than their uncertainties, resulting in a final sample size of 123 bursts. The left panel of Figure 5 shows the correlation of temperatures for the low and high temperature BB; the average values (and widths) of the log-normal distributions for the cool and hot BB are 4.6 (0.7) and 14.8 (2.1) keV, respectively.

Figure 5 (middle and right panels) also shows the correlations between the fluences and emission areas of both BBs. We estimated the significance of these correlations with a Spearman Rank Order Correlation test. In the following we designate a correlation to be significant if the chance probability, $P < 5.7 \times 10^{-7}$, corresponding to 5σ , when the correlation coefficient follows a normal distribution; a correlation is marginally significant if $5.7 \times 10^{-7} < P < 2.7 \times 10^{-3}$ (3σ). While the temperature correlation is almost significant at the 5σ level, with $P = 7 \times 10^{-7}$, the fluences and emission areas are very strongly correlated. When fitted by a power law, the slopes of the correlations are 1.86 ± 0.09 , 0.83 ± 0.02 and 1.34 ± 0.04 , for all three panels, respectively. We further discuss these correlations in Section 5.

5. DISCUSSION

We discuss here the energetics of the SGR J1550 – 5418 bursts using our spectral fitting results. We describe the correlations between the spectral and temporal parameters, and compare our results to the ones of other sources in the literature.

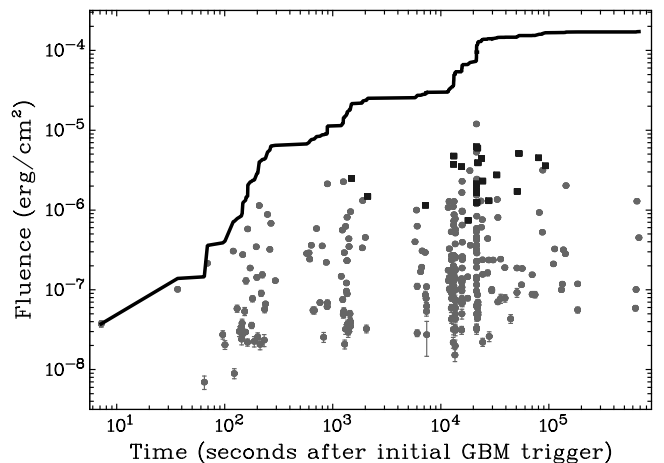


FIG. 6.— Cumulative distribution (solid line) of the energy fluence for all 286 bursts in our sample. The grey data points are individual burst fluences and the black squares indicate bursts that are partially affected by saturation.

Finally, we expand on the interpretation of our BB+BB model results.

5.1. Burst Energetics

We have determined the energy fluence for all the bursts in our sample in the 8–200 keV energy range. The fluence values obtained with our three preferred models (OTTB, Comptonized and BB+BB) are very similar. In the following we present the fluences for the Comptonized fits only, because the model has one less parameter than BB+BB, and there are several events with power-law indices deviating from -1 (corresponding to OTTB). Figure 6 shows the cumulative energy fluence during the source burst active period as well as the individual burst fluences. To obtain the best estimate of the total fluence released during that period in bursts, we have used data from all 286 bursts, including saturated events (black squares) whose values are merely lower limits. The fluence range spanned in our sample is rather large, $\sim 10^{-8}$ to $\sim 10^{-5}$ erg/cm², which is comparable to the range of SGR J0501+4516 bursts observed with GBM (Lin et al. 2011) and larger than previous studies of other magnetars (Göğüş et al. 2001; Gavriil et al. 2004). This fluence range corresponds to an energy range in bursts of $\sim 3 \times 10^{37} d_5^2$ to $\sim 3 \times 10^{40} d_5^2$ erg, which is comparable to the values found for other SGRs, but orders of magnitude higher than AXP 1E 2259+586 (Gavriil et al.

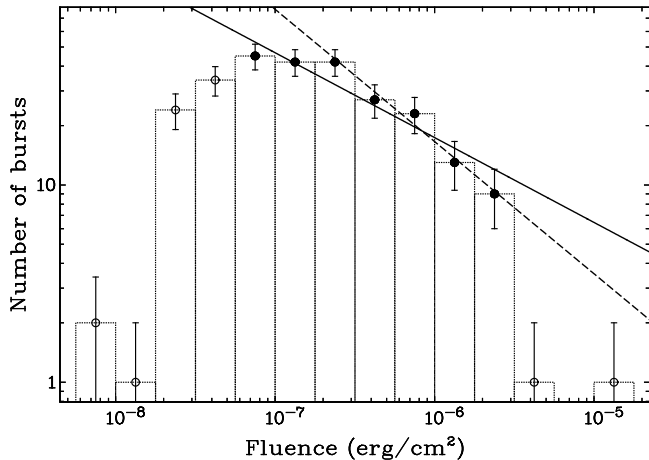


FIG. 7.— Differential distribution of the energy fluence of SGR J1550 – 5418 bursts on a logarithmic scale. The solid line shows the best power-law fit to the solid circles, ranging from $10^{-7.25}$ to $10^{-5.50}$ erg/cm², with an index of -0.4 ± 0.1 . The dashed line shows the fit of fluences ranging from $10^{-6.75}$ to $10^{-5.50}$ erg/cm², with a power-law index of -0.7 ± 0.2 . The histogram error bars indicate \sqrt{N} with N the number of bursts per fluence histogram bin. Note here that the lower fluence turn-over may be reflecting either instrumental sensitivity or an intrinsic source property, or a combination of both. Therefore we give a range for the power-law slopes.

2004).

The total energy fluence in bursts recorded with GBM alone is 1.7×10^{-4} erg/cm², which is in fact a lower limit given that there were saturated bursts, as well as bursts we excluded from our analysis since no TTE data were available, or bursts that GBM did not detect because of the SAA or Earth occultation. This results in a lower limit on the energy emitted in bursts of $5 \times 10^{41} d_5^2$ erg, which is four orders of magnitude larger than the total energy in bursts from AXP 1E2259+586 in June 2002 (Gavriil et al. 2004). The SGR J1550 – 5418 bursts detected with *INTEGRAL*/SPI-ACS on 22 January 2009 had an estimated cumulative fluence of 5.2×10^{-4} erg/cm² (25 keV–2 MeV Mereghetti et al. 2009b), i.e., a cumulative energy of $1.6 \times 10^{42} d_5^2$ erg. The total energy released in bursts detected with the *Swift* X-ray Telescope (XRT) was much lower, namely $1.6 \times 10^{40} d_5^2$ erg (1 – 10 keV; Scholz & Kaspi 2011). Scholz & Kaspi (2011) argue that the energy released in bursts is smaller than the energy released in the persistent emission between January and September 2009 over that same energy range, i.e., $1.5 \times 10^{42} d_5^2$ erg. They further claim that this behavior is similar to the one observed in AXP 1E2259+586 and not in agreement with an SGR persistent emission. Although the persistent emission energy release cannot be directly measured by GBM or SPI-ACS, Bernardini et al. (2011) have shown that for that same time period, it can be at most a factor of five higher in the 13 – 200 keV than in the 1 – 10 keV range. The energy released in bursts during a week is thus comparable to the energy released in the persistent emission during an eight month period (Scholz & Kaspi 2011), and we conclude that SGR J1550 – 5418 exhibits similar energetics behavior as other SGR sources (Woods et al. 2004).

We plot in Figure 7 the differential fluence distribution

TABLE 2
SPEARMAN RANK ORDER CORRELATION COEFFICIENTS AND PROBABILITIES FOR VARIOUS SPECTRAL AND TEMPORAL PARAMETERS

Parameters	Coefficient	Probability
$E_{\text{peak}} - \text{Fluence}$ ^a	-0.444	3.7×10^{-14}
$T_{90} - \text{Fluence}$ ^b	0.229	1.8×10^{-4}
$T_{50} - \text{Fluence}$	0.121	4.9×10^{-2}
$\tau_{90} - \text{Fluence}$ ^a	0.470	6.0×10^{-16}
$\tau_{50} - \text{Fluence}$ ^b	0.236	1.1×10^{-4}
$\delta_{90} - \text{Fluence}$ ^a	0.391	4.6×10^{-11}
$\delta_{50} - \text{Fluence}$ ^b	0.193	1.6×10^{-3}
$T_{90} - E_{\text{peak}}$	-0.071	0.25
$T_{50} - E_{\text{peak}}$	0.029	0.64
$\tau_{90} - E_{\text{peak}}$ ^b	-0.209	6.5×10^{-4}
$\tau_{50} - E_{\text{peak}}$	-0.086	0.16
$\delta_{90} - E_{\text{peak}}$ ^b	-0.252	3.4×10^{-5}
$\delta_{50} - E_{\text{peak}}$ ^b	-0.260	1.9×10^{-5}

^aSignificant correlation

^bMarginally significant correlation

for the 263 unsaturated bursts. The distribution is fitted, using unbinned maximum likelihood, between $10^{-7.25}$ and $10^{-5.50}$ with a power-law of index -0.4 ± 0.1 , and between $10^{-6.75}$ and $10^{-5.50}$ with an index -0.7 ± 0.2 . This corresponds to $dN/dF \propto F^{-1.4}$ and $dN/dF \propto F^{-1.7}$, respectively. The turn-over in the distribution at the low-fluence end is most likely caused by instrumental sensitivity drop-off; the onset of the turn-over is not well-determined and it depends on the criteria of the untriggered burst search algorithm. The high-fluence end of the distribution reflects that our sample is not complete with high-fluence bursts. When only taking the fluences above $10^{-6.75}$ erg/cm², the power-law index is comparable to the values of ~ -0.7 found for some other magnetar sources (Woods & Thompson 2006; Gavriil et al. 2004). The shallower value we obtain by including fluences down to $10^{-7.25}$ erg/cm² is closer to the values for SGR 1806-20 (Gögüş et al. 2000) and SGR J0501+4516 bursts (Lin et al. 2011).

5.2. Spectral and Temporal Correlations

Several authors have investigated correlations between the various spectral and temporal parameters of SGR bursts (e.g. Gögüş et al. 2001; Gavriil et al. 2004). Here we again present analysis using only the results of our Comptonized model fits. Using the Spearman Rank Order Correlation test we have searched for possible correlations between E_{peak} , index, fluence, and the temporal parameters $T_{90,50}$, $\tau_{90,50}$ and $\delta_{90,50}$. There are no significant correlations between the index and any other parameter; for all the other correlations the Spearman Rank Correlation coefficients and chance probabilities are listed in Table 2. From Table 2 it follows that the T_{90} duration is marginally correlated with fluence. The emission time τ_{90} is strongly correlated with fluence and the same is true for δ_{90} , while the correlations for τ_{50} and δ_{50} are only marginal. Figure 8 shows the correlations of T_{90} , τ_{90} and δ_{90} versus fluence; it is apparent that there is a stronger correlation for τ_{90} than for T_{90} or δ_{90} . When fitting a power law to these correlations, we obtain indices of 0.17 ± 0.04 for T_{90} , 0.28 ± 0.03 for τ_{90} , and 0.11 ± 0.02 for δ_{90} .

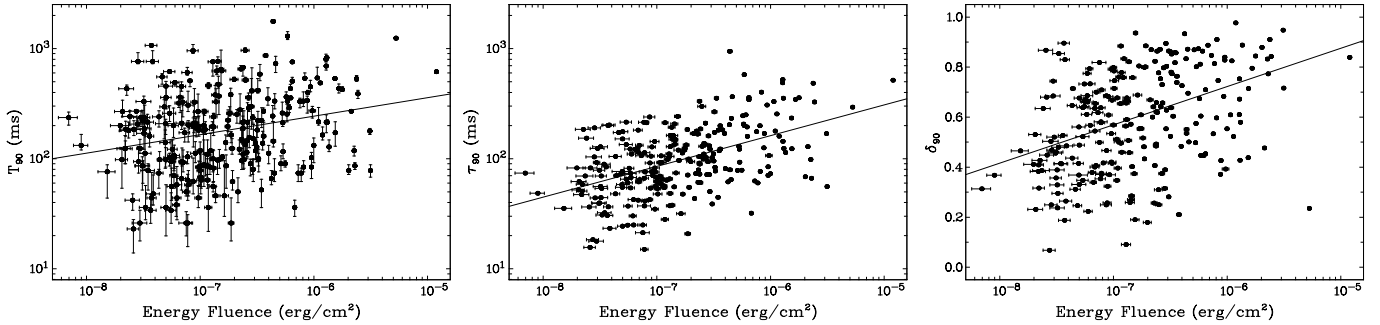


FIG. 8.— Correlation plots of T_{90} , τ_{90} and δ_{90} versus energy fluence. The lines indicate power-law fits to the data.

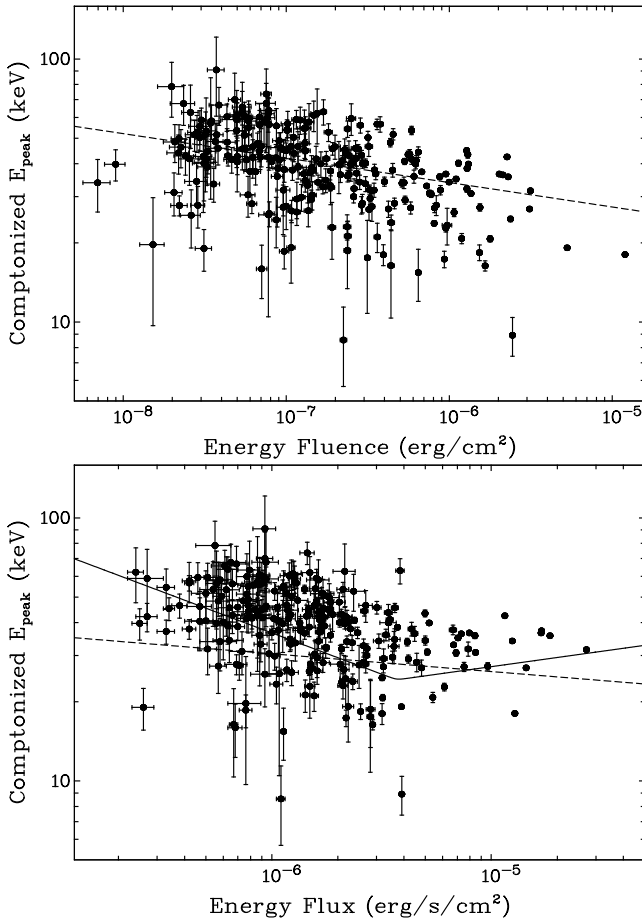


FIG. 9.— Correlation plot between Comptonized E_{peak} and energy fluence (top panel) and average energy flux (bottom panel). The dashed lines indicate a power-law fit, while the solid line in the bottom panel indicates the result of a broken power-law fit.

There is no significant correlation between E_{peak} and the duration or emission time, and there is a marginal anti-correlation between E_{peak} and the duty cycle. E_{peak} , however, is strongly anti-correlated with fluence, which is also clear from the top panel of Figure 9. The bottom panel of that Figure shows a significant anti-correlation between E_{peak} and the average burst flux, with a Spearman Rank Correlation coefficient of -0.344 and a chance probability of 9.4×10^{-9} . Since the Comptonized index does not depend significantly on the fluence of the bursts, E_{peak} is a good indicator of the

hardness of bursts. Therefore, we can conclude that there is an anti-correlation between hardness and fluence, and between hardness and average flux, for the SGR J1550 – 5418 bursts. When we fit E_{peak} versus fluence with a power law, we obtain an index of -0.093 , while an E_{peak} versus average flux fit gives an index of -0.042 (dashed lines in both panels of Figure 9).

The bottom panel of Figure 9 indicates that there may be a more complicated trend for E_{peak} versus average flux than a single power law. Similar to what has been found for the time-resolved spectroscopy of SGR J0501+4516 bursts detected with GBM (Lin et al. 2011), we find an anti-correlation at low flux values, but a positive correlation at high flux values. We fit this trend with a broken power law with indices of -0.22 and 0.07 , with the turning point occurring at an E_{peak} of 30 keV and at an average flux of 4×10^{-6} erg/s/cm². An F-test gives an improvement of 4σ (chance probability of 8×10^{-5}) for a broken power law compared to a single one. We note that a Spearman Rank Order Correlation test for the fluxes below 4×10^{-6} erg/s/cm² gives a coefficient of -0.256 and a chance probability of 7.2×10^{-5} , i.e. $\sim 4\sigma$ significance, while for fluxes above 4×10^{-6} erg/s/cm² there is not a significant positive correlation with a coefficient of 0.085 and probability of 0.64.

The minimum average flux defined by the broken power-law fit is a factor of two lower than the corresponding flux found for SGR J0501+4516 by Lin et al. (2011). Given that SGR J0501+4516 has an estimated distance of ~ 2 kpc, the minimum of an E_{peak} -luminosity correlation is three times larger for SGR J1550 – 5418 than for SGR J0501+4516. However, given the uncertainties in the source distances and in the determination of the minima in the correlations, these differences are not significant. Furthermore, these correlations were obtained from time-resolved spectroscopy in the case of SGR J0501+4516, while we present here only time-integrated results for SGR J1550 – 5418; time-resolved analysis of the brightest SGR J1550 – 5418 bursts is part of a future study. Note that in contrast with Lin et al. (2011), in the E_{peak} -fluence correlation a broken power law does not give an improvement over a single power law.

The correlation or anti-correlation between hardness and brightness of magnetar bursts has been discussed by several authors (e.g. Göğüş et al. 2001; Gavril et al. 2004; Götz et al. 2004). The spectral range and temporal capabilities of GBM allow us to characterize the burst spectral hardness with E_{peak} , as was also done in Lin et al. (2011).

This hardness indicator is much better defined than hardness ratios or indices of single power-law spectral fits. Further, our bursts span a larger fluence and flux range than any other studies with comparable sample size. Using large samples of bursts from SGRs 1806-20 and 1900+14 observed with *RXTE*, Göğüş et al. (2001) have shown that there is an anti-correlation between hardness and brightness based on hardness ratios between 10 – 60 and 2 – 10 keV. The same anti-correlation was found by Götz et al. (2004) who performed a similar analysis on *INTEGRAL*/IBIS data for SGR 1806-20 using hardness ratios between 40 – 100 and 15 – 40 keV. Contrary to the above results, Gavriil et al. (2004) showed a positive correlation for bursts of AXP 1E2259+586 detected with *RXTE*, also using hardness ratios, and concluded that this behavior distinguishes AXPs from SGRs.

For SGR J1550 – 5418, Savchenko et al. (2010) found a positive correlation between hardness and brightness based on *INTEGRAL*/SPI-ACS hardness ratios between 60 – 200 and 20 – 60 keV. However, Scholz & Kaspi (2011) do not find a hardness-fluence correlation in the *Swift*/XRT bursts on 22 January 2009. They then fit a single power law to the burst spectra, and find a significant anti-correlation between the index and flux, indicating a positive correlation between hardness and brightness. Based on this correlation, Scholz & Kaspi (2011) have argued for an AXP nature of SGR J1550 – 5418. We discuss below why the results from Savchenko et al. (2010) and Scholz & Kaspi (2011) are in apparent disagreement with our E_{peak} -fluence and E_{peak} -flux anti-correlations. The energy ranges used by Savchenko et al. (2010) and Scholz & Kaspi (2011) bracket the energy ranges used by Göğüş et al. (2001) and Gavriil et al. (2004). It is not trivial to compare hardness ratios over different energy ranges, given that the SGR burst peak emission is in the low energy band for Savchenko et al. (2010) and in the high energy band for Göğüş et al. (2001) and Gavriil et al. (2004). More importantly, we have shown here that a single power law does not provide a good description for the SGR J1550 – 5418 burst spectra, which makes the method of Scholz & Kaspi (2011) less effective than using E_{peak} as a hardness indicator. Taking all these arguments into account, we conclude that the GBM hardness-fluence anti-correlation of SGR J1550 – 5418 strongly indicates that the source is very similar to other SGR sources.

5.3. Comparison with Other SGR Sources

5.3.1. OTTB and Comptonized Model

Since OTTB or Comptonized models have very often been used as the best models for SGR bursts, we now compare our spectral fit parameters with those from studies of other SGR sources. Our E_{peak} distributions (Figure 4) are very typical for SGR bursts, but there seems to be a difference across SGR sources regarding the power-law indices of the Comptonized function. The average power-law index for SGR J1550 – 5418 bursts (-0.95) is significantly different than the one found by Lin et al. (2011) for a sample of 29 bursts detected by GBM from SGR J0501+4516, namely -0.32 . Although the width of the index distribution is fairly large for SGR J0501+4516, 0.9 compared to 0.41 for SGR J1550 – 5418, Lin et al. (2011) show that their

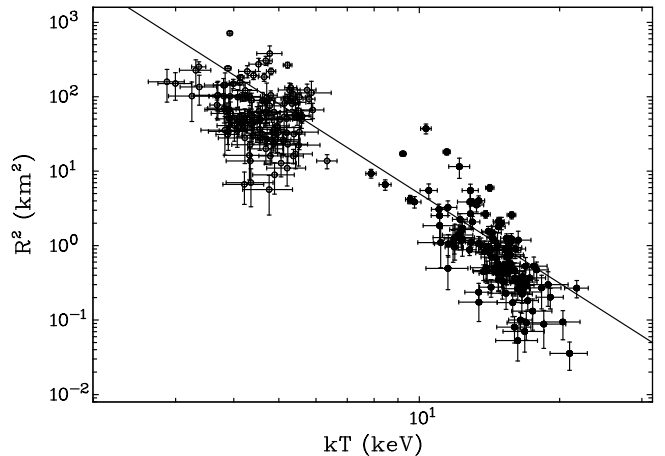


FIG. 10.— Emission area as a function of blackbody temperature for those bursts in our sample that have well constrained parameters in a BB+BB fit. The solid line indicates $R^2 \propto (kT)^{-4}$ (see text for a discussion).

brightest bursts have well-constrained indices close to 0, significantly deviating from -1 . Feroci et al. (2004) have fit 10 short bursts detected by *BeppoSAX* from SGR 1900+14 with a Comptonized function, and they find a mean spectral index of -0.4 with a dispersion of 0.2. Their result is comparable to the SGR J0501+4516 bursts and not consistent with our values for SGR J1550 – 5418. The spectral index differences between various SGR sources could be due to differences in magnetic field strength, geometry, plasma temperature or opacity. Lin et al. (2011) discuss in detail some of the effects the extreme magnetic fields of magnetars have on the Comptonized power-law index. More detailed studies, both theoretically and observationally, are needed to reach conclusive results.

5.3.2. Two blackbody Model

The temperatures, fluences and areas we find for the BB+BB model are similar to those of other SGR sources (e.g., Feroci et al. 2004; Olive et al. 2004; Nakagawa et al. 2007; Israel et al. 2008; Esposito et al. 2008; Lin et al. 2011). We have found a positive correlation between the emission areas of the two black bodies (Figure 5), which is more significant than the one found for bursts from SGRs 1806-20 and 1900+14 (Nakagawa et al. 2007). A strong correlation between the two BB luminosities has been found for SGR 1900+14 bursts with *Swift* (Israel et al. 2008) and for SGR J0501+4516 bursts with GBM (Lin et al. 2011). The latter two studies, however, have used time-resolved spectral analyses, while we show here a strong fluence correlation between the two BBs for time-integrated spectra.

Figure 10 displays the emission area as a function of temperature for both BBs. The temperature of the cool BB does not show any correlation with the fluence or the emission area of that BB. The hot BB does have a marginal anti-correlation between temperature and fluence, with a chance probability of 1×10^{-3} in a Spearman Rank Correlation test, and a very strong anti-correlation between the temperature and the emission area (Figure 10) with a chance probability of 3×10^{-24} . For illustrative purposes we have drawn a solid line in Fig-

ure 10 indicating $R^2 \propto (kT)^{-4}$, which corresponds to a constant burst luminosity or fluence. The line seems to represent best the low-temperature part of the hot BB data, while there is a clear steepening at the higher temperatures. Indeed, a power-law fit of the hot BB data gives a slope of ~ -7 , similar to what has been shown for SGR J0501+4516 by Lin et al. (2011). Similar studies have been performed for SGRs 1900+14 (Israel et al. 2008) and 1627-41 (Esposito et al. 2008) using time-resolved spectroscopy; our results confirm the trends described in these studies. A detailed time-resolved spectroscopy of SGR J1550 – 5418 bursts is underway.

Finally, we note that the smallest emission areas of the hot BB (few hundredths of km^2) are comparable to the emission area of the BB component found during enhanced persistent emission from SGR J1550 – 5418 in the GBM data after the bursting onset on 22 January 2009 (Kaneko et al. 2010), possibly indicating a common origin of the two phenomena. For the outburst data studied here, the compact hot blackbody component can be as small as $R \sim 0.2 - 0.3$ km in size, corresponding to a diameter of the order of the thickness of the outer crust, yet the bulk of the R -values for the hot BB component correspond to a diameter of ~ 1 km – comparable to the total thickness of the solid crust. For energy injection at or just below the surface, if these scales signify the size of the region in the crust that is fractured by the magnetic stresses, as is likely, then injection scales considerably smaller than the crust height have greater difficulty in disrupting it and initiating a flare, particularly if the anchored magnetic flux tube is non-radial at its footpoint. To leading order, the field energy that threads the injection region scales as R^2 , while the length of the required fracture and the corresponding energy required to produce it scale as R . Hence, crustal disruption by field twists or shear is expected to be less effective for smaller R , a contention that appears to be borne out in the data: the majority of bursts possess R values for the hot component on the scale of the thickness of the crust. Moreover, flares with a larger R for the hot BB component tend to have a larger total energy output, i.e., luminosity or fluence.

5.4. Interpretation of Two blackbody Model Results

The BB+BB model results point towards a photospheric interpretation, and perhaps one with dynamic evolution, from one magnetospheric locale to another. The fitted temperatures exceed those found (~ 0.5 keV) in the classic X-ray band (0.5 – 10 keV) in quiescent magnetar emission (e.g., see Perna et al. 2001), indicating a magnetospheric origin for such photospheres, as opposed to generation of emission purely on the neutron star surface. While it is possible that the site of outburst activation could be near the magnetic poles, the T_{90} durations of SGR J1550 – 5418 far exceed the neutron star light crossing time. Hence, even the intense magnetic fields of magnetars cannot restrict the emission region of powerful outbursts to small volumes. Accordingly, one expects spatial transport of the burst luminosity during each flare, even if tied to closed field lines. This is essentially the picture that Duncan & Thompson (1992) originally envisaged for SGRs. Polar origins would permit rapid plasma expansion along open field lines to very high altitudes, so that plasma containment is limited.

Near the equator, trapping of the gas is optimal, enabling longer durations of emission, so that quasi-equatorial locales for energy injection might be favored. The trapping times depend intimately upon the complicated interplay of polarization-dependent Compton scattering, and associated lepton diffusion and energy exchange within the photon-electron gas.

The core property evinced in Figure 10 is that the burst luminosity $\propto (kT)^4 R^2$ of the collection of bursts (and approximately also for individual bursts) is similar for the low and high temperature components. This is a total energy equipartition feature: if the surface area of a putative magnetospheric emission volume scales approximately as R^2 , then the energy in each black-body component is approximately equal. An R^2 scaling most likely applies to the hotter blackbody component as it evolves in a coronal structure, with photon and particle propagation away from an injection zone. It could be pertinent for burst activation locales near the magnetic poles, but also if the injection site is somewhat remote from the surface, for example in the magnetospheric twist scenario of Thompson & Beloborodov (2005) and Beloborodov & Thompson (2007), which was primarily envisaged for the lower luminosity, quiescent magnetar emission. However, it is possible that other volume scalings are operating, particularly due to the anisotropizing influence of the non-uniform magnetic field. For example, volumetric constriction of hot plasma near field line footpoints might modify this correlation modestly. Determining the effective volume scaling requires detailed time-dependent modeling of particle and radiative transfer in a magnetospheric environment. This motivates future theoretical work on magnetar “fireballs.” Yet, for the present, a baseline conclusion is that it is clear that the strong correlation between the energies in the two components suggests that they share the same physical origin if BB+BB is indeed the correct model for these bursts.

It is natural to expect that the high temperature component would be more closely connected to a smaller injection region, where energy is dissipated from crustal stress fractures (Duncan & Thompson 1992; Thompson & Duncan 1995) or some other origin. Proximity of the injection site to the surface is favored in such a scenario, and the resultant fireball of electrons and photons should subsequently expand and cool at higher altitudes in the closed magnetospheric zone, thereby generating the lower temperature component. The injection temperature is controlled by the total energy dissipated by hot electrons present in the inner magnetosphere. The magnetic Thomson optical depth should be high, so that Comptonization will drive thermalization, and the fireball evolution and duration will be influenced by the anisotropic, magnetic Eddington luminosity. Equilibration between photons and electrons will be non-uniform over the coronal volume, so that there should be a modest temperature gradient throughout; the emergent continuum will be a superposition of distorted blackbodies. In particular, the strong polarization dependence of the anisotropic Compton process naturally establishes two different physical scales for the optically thick environment (Ulmer 1994; Miller 1995), with the so-called E-mode photosphere being smaller than the O-mode photosphere. Perhaps these physically-distinct regions assume

somewhat different temperatures on average, though a more quantitative development of this picture is beyond the scope of this paper. Whatever the site of the injection, one anticipates that there will be distinctive evolutionary signatures in the dynamic fireball as it either expands to higher altitudes, or contracts when constricted near magnetic footpoints, resulting into a certain temporal evolution for the temperatures. Evolutionary aspects of the spectroscopy of SGR J1550 – 5418 bursts will be explored in a future paper.

6. CONCLUSIONS

We have presented the results of a time-integrated spectral and temporal study of bursts from SGR J1550 – 5418 during an extremely active bursting period, making optimal use of the spectral and temporal capabilities of GBM. The durations, emission times and duty cycles are typical of SGR bursts, with burst rise times shorter than their decays. Our spectral analysis shows that the spectra of SGR J1550 – 5418 bursts are equally well described by OTTB, Comptonized and BB+BB. The E_{peak} distributions of OTTB and Comptonized have mean values of ~ 39 keV, while the Comptonized power-law index average is close to -1 , i.e., we recover the OTTB index in our Comptonized fits. Whether this power law extends down to lower energies can only be determined by combined fits with other instruments, for instance *Swift*/XRT (Lin et al., in preparation).

We have investigated the presence of correlations between the various spectral and temporal parameters, and find significant correlations between the emission time and the fluence, and between the duty cycle and the fluence, and a marginal correlation between the duration and the fluence. We have also shown that there is a significant anti-correlation between E_{peak} and flu-

ence, and between E_{peak} and average flux. We find an anti-correlation at low flux levels but a weak correlation with E_{peak} at high fluxes, with a possible minimum at 4 erg/s/cm^2 . In a time-resolved analysis of the bursts from SGR 0501+4516, the positive correlation behavior at high fluxes is more significant. Whether the two sources have a similar behavior remains to be seen in the time-resolved analysis of the SGR J1550 – 5418 bursts. One of the main contributions of the GBM data is providing accurate measurements of the evolution of E_{peak} for short bursts, which is the best hardness indicator for hardness-brightness correlation studies.

For the BB+BB fits we find average temperatures of ~ 5 and ~ 14 keV and we show that these temperatures are well correlated. The fluences of the two BB functions are strongly correlated, and the same is true for their emission areas. The emission area of the low-temperature BB is comparable to the neutron star surface area; the hot BB area is much smaller and strongly anti-correlated with its temperature.

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